

1. Goodman 6-2: The *line-spread function* of a two-dimensional imaging system is defined to be the response of that system to a one-dimensional Dirac delta function passing through the origin of the input plane.

- (a) In the case of a line excitation lying along the x -axis, show that the line-spread function $\ell[y]$ and the point spread function $p[x, y]$ are related by:

$$\ell[y] = \int_{-\infty}^{+\infty} p[x, y] dx$$

where ℓ and p are interpreted as amplitudes or intensities, depending on whether the system is coherent or incoherent, respectively.

- (b) Show that for a line source oriented along the x -axis, the 1-D Fourier transform of the line-spread function is equal to a slice through the 2D Fourier transform of the point-spread function along the ξ -axis. In other words, if $\mathcal{F}_1\{\ell[y]\} = L[\eta]$ and $\mathcal{F}_2\{p[x, y]\} = P[\xi, \eta]$, then $L[\eta] = P[0, \eta]$ (this is the so-called *central-slice theorem*).
- (c) Find the relationship between the lsf and the step response of the system, i.e., the response to a unit-step excitation oriented parallel to the x -axis.
- (d) (ADDITIONAL PART) This method may be adapted to evaluate the OTF of an imaging system with a nonsymmetric pupil, as from systems composed of multiple apertures. Describe how this may be accomplished.
2. Goodman 6-6:
- (a) Sketch the cross sections of the OTF of an incoherent imaging system whose pupil consists of two circular apertures of diameter $2w$ separated by $2d$. Label the various cutoff and center frequencies.
- (b) (ADDITIONAL PART) Repeat (a) for the case of an incoherent system where one of the two “subapertures” is misaligned in its cell to add a phase ϕ to the light through that aperture. For example, if the system is a reflector telescope, one of the two mirrors can “slip” in its cell by some amount Δz so that the optical path difference is $\phi = \frac{2\pi}{\lambda} \cdot \Delta z$. Evaluate for the cases $\phi = \frac{\pi}{4}$, $\frac{\pi}{2}$, and π radians.

3. Goodman 6-8: Consider the incoherent OTF for focusing error in Eq.(6-41):

$$\mathcal{H}[\xi, \eta] = TRI\left[\frac{\xi}{2f_o}, \frac{\eta}{2f_o}\right] SINC\left[\frac{8W}{\lambda_0}\left(\frac{\xi}{2f_o}\right)\left(1 - \frac{|\xi|}{2f_o}\right)\right] \\ \times SINC\left[\frac{8W}{\lambda_0}\left(\frac{\eta}{2f_o}\right)\left(1 - \frac{|\eta|}{2f_o}\right)\right] \quad (6-41)$$

as predicted for a system having a square pupil and focusing error. It is hypothesized that the psf of this system is the convolution of the diffraction-limited psf with the psf predicted by geometrical optics. Examine the validity of this claim.

4. Goodman 6-9: A metric of considerable utility in determining the “seriousness” of aberrations of an optical system is the *Strehl definition* \mathcal{D} , which is defined as the ratio of the light intensity at the maximum of the psf of the system with aberrations to that same maximum for an aberration-free system. (Both maxima are assumed to exist on the optical axis). Prove that \mathcal{D} is equal to the normalized volume under the OTF of the aberrated system; that is, prove

$$\mathcal{D} = \frac{\iint_{-\infty}^{+\infty} \mathcal{H}[\xi, \eta]_{\text{with aberrations}} d\xi d\theta}{\iint_{-\infty}^{+\infty} \mathcal{H}[\xi, \eta]_{\text{without aberrations}} d\xi d\theta}$$

5. Goodman 6-13: The *F-number* of a lens with a circular aperture is defined as the ratio of the focal length to the lens diameter. Show that thwn the object distance is infinite, the cutoff frequency for a coherent imaging system using this lens is given by

$$f_o = \frac{1}{2 \lambda F/\#}$$